Numerical calculations of the thermal-aerodynamic characteristics in a solar duct with multiple V-baffles

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ABSTRACT

The study aimed to enhance the heat transport by improving the hydrodynamic structure of the system by changing and restructuring the duct’s internal geometry. Modern fins, of the shape ‘V’, have been proposed with different dimensions, and they are periodically arranged over the duct surfaces. The most important steps of this research are the change in the V-fin attack-angle (40°–80°), length (Hb/2, 3Hb/4, Hb, 5Hb/4 and 3Hb/2), and separation length (Ds/2, 3Ds/4, Ds and 5Ds/4), as well as the flow rate (6×10³–3×10⁴). The study yielded an optimum case for a 40-degree attack-angle, with a factor of thermal enhancement of 2.163 for the highest value of Reynolds number. On the other hand, improving the length of the V-fins or decreasing in the space between them, increases the flow strength by enlarging the recycling cells, which reflects on the hydrodynamic behavior, and changes the heat transfer. The presence of this new model of fins also highlights a hydrothermal improvement ranging between 1.196 and 23.779 percent compared to the previously indicated models, reflecting the effectiveness of the new system of solar heat exchangers with air V-finned ducts.

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KEYWORDS

CFD; solar heat exchangers; air V-finned ducts; recycling cells; heat transfer

Nomenclature

\[ \text{Nu} \] \quad \text{average number of Nusselt}  
\[ \text{Nu}_0 \] \quad \text{Nusselt number for smooth exchanger}  
\[ \text{Nu}_x \] \quad \text{local number of Nusselt}  
\[ p \] \quad \text{pressure [Pa]}  
\[ P_d \] \quad \text{dynamic pressure [Pa]}  
\[ Pr \] \quad \text{number of Prandtl}  
\[ Re \] \quad \text{number of Reynolds}  
\[ T \] \quad \text{temperature [K]}  
\[ u \] \quad \text{horizontal velocity [m/s]}  
\[ U_{in} \] \quad \text{intake velocity [m/s]}  
\[ v \] \quad \text{vertical velocity [m/s]}  
\[ V \] \quad \text{mean velocity [m/s]}  
\[ w \] \quad \text{thickness of fin [m]}  
\[ W \] \quad \text{width of exchanger [m]}  
\[ \Delta P \] \quad \text{pressure drop [Pa]}  

Greek letter

\[ TEF \] \quad \text{factor of thermal performance}  
\[ \varepsilon \] \quad \text{rate of dissipation of turbulence [m}^2/\text{s}^3\text{]}
\( \alpha \) angle of attack of fin [degree]
\( \mu \) dynamic viscosity [kg/m s]
\( \mu_t \) turbulent viscosity [kg/m s]
\( \rho \) density of fluid [kg/m\(^3\)]
\( \sigma_k \) constant in \( K \)-equation
\( \sigma_\varepsilon \) constant in \( \varepsilon \)-equation
\( \tau_w \) wall-shear-stress [Pa]
\( \psi_r \) stream-function [kg/s]

**Subscript**

0 smooth
a aeraulic
b baffle
c exchanger conduit
f fluid
in intake
max maximum
out outlet
ref reference
s separation
t turbulent
x local

**Introduction**

Solar energy is modern, pure, renewable, and inexhaustible (Ahmadi et al., 2020; Beyaztas et al., 2019; Samadianfard et al., 2019). Transforming solar energy into thermal energy via solar collectors is a topic of interest to many researchers (Abuşka, 2018; Kabeel et al., 2019). Solar air duct heat exchangers target many studies (Amraoui & Aliane, 2018; Hassan & Abo-Elfadl, 2018; Ho et al., 2017). The inclusion of obstacle-type vortex generators is one of the most effective and ongoing techniques to date (Ameur, 2020; Ameur & Menni, 2019; Jain et al., 2019; Kumar et al., 2018; Kumar & Kim, 2015). We present in this axis some of the numerical and experimental studies that have become popular in this field, starting with researchers Promvonge and Thianpong (2008). Some experiments have been carried out by these authors for evaluating friction loss and turbulent forced-convective heat transfer of airflow within a steady heat flux channel equipped with different-configurations ribs. The structures of the ribs were rectangular, wedge, and triangular. The study showed an increase in the value of Nusselt number as well as in the coefficient of skin friction in the case of a wedge-section rib, while the highest thermal performance was recorded in the case of a triangular-geometry rib with a staggered arrangement. Furthermore, convective transfer of heat, as well as laminar fluid flow attributes in a 2-D horizontal channel with adding diamond-shaped baffles under the condition of isothermal walls have been numerically studied by Sripattanapipat and Promvonge (2009). In this study, outcomes of distinct baffle tip angles on pressure loss as well as thermal performance within the channel have been evaluated, and comparisons of using flat and diamond baffles have been provided. A comparison of different diamond baffles showed high thermal performance for the attack angle of 5 degrees, while the flat baffle showed a decrease in the same factor and for all Reynolds values followed, compared to the diamond baffle with an angle of attack varying from 5 to 10 degrees. In another study, hydrodynamic behavior was detected as well as an examination of heat transfer in a rectangular structure of a channel equipped with pin-typed fins (Wang et al., 2012). Also, the comparison of various shaped pin fins with identical cross-sectional areas such as drop-shaped, elliptical, and circular in a staggered arrangement has been presented. The results showed an improvement in the thermal performance in the case of circular-type pins compared to the case of drop-geometry pins.

Furthermore, Reddy and Satyanarayana (2008) performed a 3-D CFD simulation of a new solar receiver for the aim of improving thermal performance. The receiver contained different porous inserts like circular, trapezoidal, and square shapes. The comparative results showed that the trapezoidal-shape porous-type inserts increase the heat transfer by about 13.8% in the case of 1.7 kPa of pressure penalty and 6.4 kg/s of fluid mass. Using the FLO EFD 9.0 type CFD software, Xu et al. (2014) conducted a detailed numerical analysis of the flow fields as well as the isotherms in a computer board. The whole system adopted the combination of liquid-type cooling and air-type cooling to increase heat transfer. The computational results showed an improvement in the heat dissipation efficiency in the presence of fin-type heat sinks, as the temperature of the CPU case dropped from 186° to 149°C. Also, Jayavel and Tiwari (2010) used a 3D CFD code to address the hydrothermal characteristics through a rectangular-configuration channel by using the tube separation technique. Their results reported that the heat transfer could be enhanced in both cases of the tube arrangement, i.e. in-line and staggered, by reducing the separation length between them. Vatani and Mohammed (2013) applied various rib-groove structures such as arc, square, and triangles with different nanofluids containing SiO\(_2\), ZnO, CuO, Al\(_2\)O\(_3\) to numerically investigate hydraulic and thermal characteristics of flow in the channel. As indicated in the results, the greatest Nusselt number can be achieved by using SiO\(_2\)/water nanofluid among all the nanofluids. Moreover, in another CFD study, different rib orientations have been used by Gao et al. (2018) with the
The aim of improving the heat transfer in a U-duct when $Re = 3 \times 10^4$. The study showed that the recirculation core technique is an important visualization way for flow checking and optimization in U-section cooling-type ducts with angled-geometry ribs. Also, Cao et al. (2019) conducted a study by using experimental techniques as well as numerical schemes on a new heat exchanger model with inserting sextant helical baffle for the aim of improving thermal performance. The analysis of comparison results revealed that the efficiency of the sextant-type scheme is greater than that of the quadrant one. In addition, the authors indicated that this study can be adopted as a theoretical database for updating heat exchangers according to new models.

Saini and Saini (2008) conducted a study by using experimental methods for developing the heat in a solar heater with airflow and arc-shape parallel wire to increase roughness. Furthermore, they investigated the friction factor augmentation caused by these elements. Their study obtained a maximum thermal improvement of about 3.8 times for the relative angle of the wire estimated at 0.3333 with a relative height of 0.0422. This improvement also corresponds to an increase in the factor of friction, as it was estimated by about 1.75 times at the same previous dimensions. Skullong et al. (2014) conducted a study by using experimental methods for investigating thermal performance and airflow characteristics into a great solar heater duct given by groove-type turbulators and wavy-type ribs. Their study showed an enhancement in the Nusselt number and friction coefficient for the case of simultaneous use of rib-groove on both duct surfaces compared to the smooth-type duct in the presence or absence of ribs. Also, an innovative method to improve the transfer of heat and decrease friction-loss inside a heat exchanger provided by circular-type tubes have been introduced by Torii et al. (2002) using turbulators of winglet-type. The results reveal that there is a 30% increase in heat transfer in the case of overlapping tube banks, while it is 20% in the inline-manner tube bank second case. In another study, an experimental comparison was made by Zhou and Ye (2012) between conventional winglet-type vortex generators, i.e. delta, rectangular and trapezoidal, with novel winglet-type vortex generators, i.e. curved trapezoidal. The results concluded with the selection of the traditional model of vortex generators with a delta-type winglet for laminar flows as well as transitional streams, while the effectiveness of the new winglet model of curved trapezoidal type was confirmed for completely turbulent flows. Additionally, an especial flow layout has been designed by Du et al. (2017) to check the molten salt performance for heat transfer in a heat exchanger fitted with fractional obstacles in Reynolds numbers between 6,142 and 9,125 and high temperature. The authors reported an increase of 26% in values of Nusselt number, in the cases of using segmental obstacles and molten salt for a low stream rate. Lei et al. (2008) simulated by using numerical approaches the effectiveness of the attack value of helical-form baffles on the performance of heat exchangers. According to the results, the mean Nusselt number is augmented by increasing the baffle inclination. Additionally, Wen et al. (2015) developed a pattern of a ladder-type baffle for the purpose of enhancing the thermal performance of heat exchangers by utilizing helical baffles. The results confirmed the importance of using the foldable ladder-model of helical baffles to enhance thermal performance in such heat exchangers used. Also, an experiment has been conducted by Bopche and Tandale (2009) to study skin friction and convective coefficient of heat transfer in the solar duct from an air heater by inserting U-shaped turbulators. The study showed an augmentation of about 2.82 times in terms of heat transfer, while 3.72 times in terms of friction, compared to the smooth-wall duct. Basing on experimental approaches, Promvonge et al. (2015) evaluated the influences of width ratios and pitch length for horseshoe-type turbulators with 20° inclination on the hydrothermal aspects of a tubular section heat exchanger. As indicated in the results, the heat transfer rate improved significantly. Skullong et al. (2016) reported a computational and applied study in order to collect the impacts of using 30° horseshoe-baffle (HB) in a square duct on thermal performance and turbulent flow. Comparing the performance of using these HBs, it turns out that there is a high improvement received from this horseshoe-type new model, much better than in the case of the presence of wire coils. Also, Singh et al. (2011) used experimental methods to determine the effects of using discrete V-down rib in a rectangular duct on the fluid flow as well as the thermal characteristics with consistent heat flux. Their study concluded to demonstrate the effectiveness of the new baffle structure, which gave an increase in Nusselt number by about 3.04 folds and an augmentation in skin friction by about 3.11 times relative to the smooth-wall duct. Boonlai and Jedsadaratanachai (2016) numerically investigated thermal performance enhancement of a square-geometry duct containing discrete-section combined-model vortex generators. In their simulation, the effects of flow directions and blockage ratios in a turbulent regime, i.e. 3,000 < $Re$ < 20,000, have been analyzed. Under the conditions used, the thermal transfer enhancement has been updated about 2.8 to 6 times if compared to the results of the smooth-walled channel. Furthermore, an experiment has been carried out by Chamoli and Thakur (2016) with the purpose of determining friction factor along with the thermal characteristics in a rectangular
Table 1. New physical domain with novel vortex generators.

<table>
<thead>
<tr>
<th>Author(s)</th>
<th>Physical model</th>
<th>Physical domain</th>
</tr>
</thead>
</table>
| Pourramezan et al. (2020) | Circular type tube with twisted-conical-strip form inserts | - Laminar flow
- Taguchi approach
- ANSYS Fluent 14.5 |
| Saravanakumar et al. (2020) | Air heater with arc-type ribs | - Extended-type surface
- Exergetic analysis
- Optimization
- MATLAB |
| Nakhchi et al. (2020) | Heat exchangers with perforated-louvered-strip type inserts | - \( Re = 5,000-14,000 \)
- RNG type k-\( \varepsilon \)
- Ansys ICEM CFD 18.1 |
| Saleh et al. (2019) | Square-form enclosure with 2 elastic-type fins | - Laminar flow
- Mixed-convection
- Arbitrary Lagrangian-Eulerian |
| Yu et al. (2020) | Shell and tube type HE with longitudinal VGs | - Turbulent flow
- RSM numerical study |
| Promvonge and Skullong (2019a) | Solar receiver with hole punched type wings | - Heat transfer
- Turbulent flow
- \( Re = 5,300-22,600 \)
- Experimental study |

(continued)
### Table 1. Continued.

<table>
<thead>
<tr>
<th>Author(s)</th>
<th>Physical model</th>
<th>Physical domain</th>
</tr>
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<tbody>
<tr>
<td>Promvonge and Skullong (2020)</td>
<td>Tubular type HE with V-baffled type tapes</td>
<td>Forced convection, $Re = 4,120–25,800$, ANSYS Fluent</td>
</tr>
<tr>
<td>Promvonge and Skullong (2019b)</td>
<td>Solar receiver with combined turbulators</td>
<td>Forced convection, Turbulent flow, $Re = 5,300–23,000$, Experimental study</td>
</tr>
<tr>
<td>Ekiciler (2020)</td>
<td>Duct with triangular-type rib</td>
<td>Hybrid nanofluid, Turbulent flow, $Re = 50,000–100,000$, Finite volume, RING type k-epsilon</td>
</tr>
<tr>
<td>Eiamsa-ard and Chuwattanakul (2020)</td>
<td>Channels with twisted type baffles</td>
<td>Heat transfer, $Re = 4,000–20,000$, Experimental study</td>
</tr>
<tr>
<td>Liang et al. (2019)</td>
<td>Tubes with inserts</td>
<td>Water type fluid, Heat transfer, Laminar flow, $Re = 6,00–1,800$, Numerical study, Finite volume, FLUENT 14.0</td>
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(continued)
duct by inserting V-down baffles. The study concerned with the treatment of enhancing the heat transfer by using the roughness structure for the duct walls, where its various effects on the hydrothermal behavior were determined, and the largest heat transfer was obtained for a relative pitch ranging between 1.5 and 3. Using the technique of boundary elements as well as the approach of finite elements, Ghalandari et al. (2019) conducted numerical studies to examine various sloshing terms as well as the flexibility terms of elastic-type structures, submerged in a flow filed. According to the obtained data, the change in the dimensions of the elastic structures, such as width, thicknesses as well as height, strongly changes the hydrodynamic behavior of the coupled system. Salih et al. (2019) conducted a detailed examination and full evaluation of the interactions between the moving limits of the thin-type structures and the flow fields by using numerical algorithms and methods. Akbarian et al. (2018) used natural-type gas for dual-fuel model engines through computational and applied studies. Additionally, Jedsadaratanachai and Boonloi (2014) and Sirromreun et al. (2012) determined flow friction along with thermal characteristic of a channel with V-shaped and Z-shaped baffles, respectively. Other theoretical and experimental experiments, such as Pourramezan et al. (2020), Saravanakumar et al. (2020), Nakhchi et al. (2020), Saleh et al. (2019), Yu et al. (2020), Promvonge and Skullong (2019a, 2019b, 2020), Ekiciler (2020), Eiamsa-ard and Chuwattanakul (2020), Liang et al. (2019), Zhang et al. (2019) and Biçer et al. (2020), considered and studied new physical models and domains with novel vortex generators as addressed below in Table 1.

According to other ways, many other studies have resorted to exploiting new fluids as a substitute for conventional fluids, i.e. air and water, as reported by Dogonchi et al. (2019, 2020); Hoseinzadeh et al. (2019); Mehryan et al. (2019); Alsabery et al. (2019); Kumar et al. (2020); Chamkha and Al-Mudhaf (2005); Chamkha et al. (2016, 2017); Zaraki et al. (2015); Goria and Chamkha (2011); and Mohebbi et al. (2017); Takhar et al. (2001) for different thermal systems.

Solar air duct heat exchangers are involved in many applications such as heating, cooling and drying. Therefore, it is important to update these thermal devices by giving them an effective engineering structure by inserting new vortex generators such as ‘V’-shaped obstacles. The finning technology has been the subject of many previous studies where fins in their traditional form have been used strongly. Also, some authors have considered other modern models, such as helical obstacles. Sure, this type of vortex generators is important in terms of performance, but it is complex in structure and difficult to achieve.

On the other hand, other new high-performance configurations have been used, according to different forms, such as V, W, and Z, where they can easily be exploited experimentally. Therefore, it is important to activate such models in raising the performance of some solar ducts

### Table 1. Continued.

<table>
<thead>
<tr>
<th>Author(s)</th>
<th>Physical model</th>
<th>Physical domain</th>
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<tr>
<td>Zhang et al. (2019)</td>
<td>- Helical type channel with VGs</td>
<td>- Experimental study</td>
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<td></td>
<td>- 3D simulation</td>
<td>- SIMPLE-algorithm</td>
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<td></td>
<td>- ANSYS FLUENT</td>
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<tr>
<td>Biçer et al. (2020)</td>
<td>- HE with three-zonal type baffles</td>
<td>- Turbulent flow</td>
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<tr>
<td></td>
<td>- Optimization</td>
<td>- Experimental study</td>
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<tr>
<td></td>
<td>- Taguchi-method</td>
<td>- Standard type k-ε</td>
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<td></td>
<td>- ANSYS Fluent</td>
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for air-heat exchangers by conducting research studies using the CFD codes. Through this research, we present a numerical simulation of thermal-aerodynamic characteristics of a solar air duct with multiple V-shaped baffles. The thermal-aerodynamic characteristics are listed below:

- The aerodynamic characteristics are: streamlines (or stream-function), mean, axial and transverse velocities, dynamic pressure, as well as local and average friction coefficients (or pressure drop).
- The thermal characteristics are: isotherms, as well as local and average Nusselt numbers.
- The thermal-aerodynamic characteristic is the thermal enhancement factor (Performance: TEF).

In this study, we focus mainly on important geometric variables related to V-fins, namely the attack angle, the height as well as the separation length between them, along with the flow rate, in order to reach an optimal solution for better conditions.

**Conception of a duct provided by vortex V-generators**

Figure 1 is a cross-section diagram of a 2-D horizontal rectangular duct model from an air-heat exchanger with novel fin-type obstacles in the upstream V-geometry in order to change the main flow path, to create recycling areas, to ensure good mixing, and to raise the transport of the heat from the hot region to working fluid.

The new heat exchanger model, reported in Figure 1, has been prepared on the basis of empirical and numerical data from Demartini et al. (2004). In their studies, the airfield was treated in the case of a pair of vertical simple obstacles having a height \( H_b \) of \( 8 \times 10^{-2} \) m and a thickness \( w \) of \( 10^{-2} \) m with a pitch \( D_s \) of \( 142 \times 10^{-3} \) m, inside a heat exchanger of \( 146 \times 10^{-3} \) m in height \( H_c \) and \( 167 \times 10^{-3} \) m in aerulic diameter \( D_h \).

This contribution aims to update this heat exchanger by introducing a new model of staggered upstream V-shaped fins as an alternative to simple obstacles. This study conducts two-dimensional simulations to evaluate the hydrothermal and dynamic structure of the multi-finned duct in the presence of different attack angle \( \alpha \), varying from 40° to 80° with respect to the horizon, and multiple heights, from \( H_b/2 \) to \( 3H_b/2 \), with variable separation lengths, from \( D_s/2 \) to \( 5D_s/4 \). Moreover, compared to other models of the ducts, such as smooth ducts or those supported by obstacles, i.e. 90° simple, downstream-type V, upstream-type V, trapezoidal, triangular, S, W and Z vortex generators, they are also highlighted and discussed in order to show the effectiveness of the current model.

The aerulic boundary condition is considered to be a consistent 1-D velocity \( U_{in} \) at the duct’s intake as provided in the literature (Demartini et al., 2004; Endres & Möller, 2001; Nasiruddin & Siddiqui, 2007). The air temperature is 300 K at section \( x = 0 \) (Nasiruddin & Siddiqui, 2007). The solid limits of the exchanger, as well as the fin installation bases, are kept to have 375 K constant temperature (Nasiruddin & Siddiqui, 2007). The conditions of wall limits are considered to be non-slip, and atmospheric pressure is applied in the exchanger outlet (Demartini et al., 2004). Additionally, the succeeding presumptions are considered to develop the hydrothermal improvement: both the air-fluid and the aluminum metal used have fixed properties, the aerulic fluid is considered Newtonian as well as incompressible, while the flow is steady-state turbulent \( (Re = 6 \times 10^3 – 3 \times 10^4) \). Also, the radiation heat transfer is neglected.

**Mathematical formulation and numerical solution**

Under the abovementioned presumptions, the hydrothermal model is controlled by (Nasiruddin & Siddiqui, 2007):

The continuity:

\[
\nabla \cdot \vec{V} = 0
\]

The momentum:

\[
\rho (\vec{V} \cdot \nabla) \vec{V} = -\nabla P + \mu f \nabla^2 \vec{V}
\]
The energy:

$$\rho Cp (\mathbf{\nabla} \cdot \mathbf{V}) = k_f \nabla^2 T$$  \hspace{1cm} (3)

The turbulence phenomenon is mathematically modeled by using the $k-\varepsilon$ model of standard-type (Launder & Spalding, 1974). It has a first equation, related to kinetic-energy of turbulence ($K$), which is defined as below:

$$\frac{\partial}{\partial x_j} (\rho k u_j) = \frac{\partial}{\partial x_j} \left[ \left( \mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + G_k + \rho \varepsilon$$  \hspace{1cm} (4)

While the rate of dissipation of turbulence ($\varepsilon$) is the content of the second equation:

$$\frac{\partial}{\partial x_j} (\rho \varepsilon u_j) = \frac{\partial}{\partial x_j} \left[ \left( \mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial \varepsilon}{\partial x_j} \right] + C_{1\varepsilon} \frac{\varepsilon}{k} - C_{2\varepsilon} \rho \frac{\varepsilon^2}{k}$$  \hspace{1cm} (5)

The aeraulic-diameter:

$$D_h = 2HW/(H+W)$$  \hspace{1cm} (6)

The Reynolds-number:

$$Re = \rho UD_h/\mu$$  \hspace{1cm} (7)

The local-Nusselt-number:

$$Nu_x = \frac{h(x)D_h}{k_f}$$  \hspace{1cm} (8)

The average-Nusselt-number:

$$Nu = \frac{1}{L} \int Nu_x \, dx$$  \hspace{1cm} (9)

The Nusselt-number correlation of Dittus and Boelter (1930):

$$Nu_0 = 0.023Re^{0.8}Pr^{0.4}$$  \hspace{1cm} (10)

The local-friction-coefficient:

$$C_f = \frac{2\tau_w}{\rho U^2}$$  \hspace{1cm} (11)

The average-friction-coefficient:

$$f = \frac{2(\Delta P/L)D_h}{\rho U^2}$$  \hspace{1cm} (12)

The friction-factor correlation of Petukhov (1970)

$$f_0 = (0.79 \ln Re - 1.64)^{-2}$$  \hspace{1cm} (13)

The performance-factor:

$$\eta = \frac{(Nu/Nu_0)}{(f/f_0)^{1/3}}$$  \hspace{1cm} (14)

Furthermore in this study, the employed mesh structure is a 2-D, non-uniform grid. The refined mesh at all solid walls and limits is shown in Figure 2. This refinement is needed to have precise temperature gradients and velocity in the area under study. As illustrated in Figure 2, the mesh far from the walls is uniform.

This simulation has been tried out using two basic components, namely the computational technique of finite volumes (Patankar, 1980), Algorithm of discretization ‘SIMPLE’ (Patankar, 1980), through the software ‘ANSYS Fluent 12.0 (2012)’. Additionally, Leonard and Mokhtari (1990)’s interpolation QUICK-scheme has been employed for the flow properties. And, for the pressure terms, a second-order upwind method (Patankar, 1980) was implemented, see Table 2.

On the other hand, the intensity of the mesh and its effect on the computational solution are investigated as shown in Table 3. The qualitative changes of maximum dynamic pressure ($P_{d_{\text{max}}}$) and x-velocity ($u_{\text{max}}$) values are analyzed by giving their relative percent errors, respectively, $\delta$ (see Eq. 15a) and $\delta'$ (see Eq. 15b), in the presence of different cell meshes.

$$\delta = \frac{P_{d_{\text{ref}}} - P_d}{P_{d_{\text{ref}}}} \times 100 \hspace{1cm} (15a)$$

$$\delta' = \frac{u_{\text{ref}} - u}{u_{\text{ref}}} \times 100 \hspace{1cm} (15b)$$

The least relative error is observed on the grid of $245 \times 95$ nodes, as it reached about 0.130% in the case

### Table 2. Solution controls.

<table>
<thead>
<tr>
<th>Discretization</th>
<th>Second-Order Upwind</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pressure</td>
<td>Quick</td>
</tr>
<tr>
<td>Momentum</td>
<td>Quick</td>
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<tr>
<td>Kinetic energy of turbulence</td>
<td>Quick</td>
</tr>
<tr>
<td>Rate of dissipation of turbulence</td>
<td>Quick</td>
</tr>
<tr>
<td>Energy</td>
<td>Quick</td>
</tr>
<tr>
<td>Pressure-velocity coupling</td>
<td>SIMPLE</td>
</tr>
</tbody>
</table>

**Figure 2.** Used non-uniform grid.
of $Pd_{\text{max}}$ while 0.244% in the case of $u_{\text{max}}$, by raising it to $370 \times 145$ cells. So, the grid of $245 \times 95$ cells is used for all simulations. Also for solution convergence, the normalized residuals are fixed at $10^{-9}$ for equations of continuity (Eq. 1), $x$- and $y$- velocities (Eq. 2); $k$ (Eq. 4) and $\varepsilon$ (Eq. 5), while $10^{-12}$ for the energy equation (Eq. 3).

In the beginning, an attempt is made to validate some predicted results against the experimental data of the literature. A reference is made to the experimental work realized by Demartini et al. (2004), which consists of a simple duct heat exchanger as that considered by Demartini et al. (2004). With the same geometrical conditions as those undertaken by these authors, the values of the pressure coefficient ($C_p$) above the second plate at $x = 0.375$ m are depicted in Figure 3a along with the exchanger height. Furthermore, the variation of axial velocity ($u$) in the upper part of the duct at $x = 0.525$ m (i.e. near the outlet section) is illustrated in Figure 3b. As clearly observed, the comparison of our numerical results with those obtained numerically and experimentally by Demartini et al. (2004) reveals a satisfactory agreement.

Another verification of the reliability of our numerical method is performed in Figure 4, where the effect of the turbulence models on some predicted results is inspected. Four turbulence models were examined for this purpose, namely: standard-type $k$-$\varepsilon$, RNG-type $k$-$\omega$, standard-type $k$-$\omega$, and SST-type $k$-$\omega$ models. The curves of axial velocity in the top region of the channel at $x = 0.525$ m and $U_{\text{in}} = 7.8$ m/s are presented in Figure 4 for the four studied models of the turbulence. These results are compared to the experimental data of Demartini et al. (2004). From this figure, it seems that the standard-type $k$-$\varepsilon$ is the model that gives the lowest deviation from the experimental data. Therefore, this model has been selected to achieve the next computations.

### Results and discussion

#### Effect of the attack angle

In this section, the effect of attack angle of the V-fins ($\alpha$) is examined. The $\alpha$ angle is changed from 40° to 90°...
with a degree step of 10°, i.e. six geometrical cases were considered for constant V-fins’ height and space values ($H_b = 0.08$ m and $D_s = 0.142$ m).

Figure 5a-g highlights the axial velocity field values for solar duct air-heat exchangers without or with vortex generators (VGs) according to different attack angles varying from 40 to 90 degrees, at $Re = 6 \times 10^3$. It is very clear that the current of air is stable in the absence of VGs. In the case of V-finned heat exchangers, the amounts of the $x$-velocity can be relatively negligible adjacent to the V-fins, specifically within the downstream zones since the recirculation areas exist. Away from these areas, the lines of the air-current turn into a parallel form that gradually engenders the flow development. It should be noticed that the $x$-velocity values increase in the area expanded from the VG end to the exchanger’s wall. The reason for the $x$-velocity increase can be mentioned as using V-fins and creating the circulating motion in the duct. Thus, a sudden alter in the flow direction appears. Also, it can be noticed that the highest amount of $x$-velocity can be observed adjacent to the exchanger top for all V-fins. Furthermore, the V-fin’s attach angle impacts the flow characteristics. If the attack angle ($\alpha$) is increased, it leads to the flow deviation and acceleration adjacent to the V-fins which results in intensifying the convective rate of heat transfer. While this attack angle augmentation is restricted since it might cause the pressure drop to augment. It seems evident which more sensible effects on the upper left side of the V-fin can be seen by altering the V-angle ($\alpha$) because of the air-field direction deviation. On the other hand, increasing the $\alpha$ angle from 80° to 90° increases the field intensity and enhances the recirculation cells, giving maximum values of $x$-velocity.

In terms of the velocity’s transverse component ($y$-velocity), it is worth stating which for all the cases examined; negative $y$-velocity gradients can be noticed at the top-surface installed V-fin tip and positive $y$-velocity gradients at the bottom-surface placed V-fin end, Figure 6a-g. The $y$-velocity intensifies by enhancing $\alpha$ angle and consequently, the 90° case (simple baffle) provides maximum $y$-velocity.

Figure 7a-g illustrates the dynamic pressure ($Pd$) field distribution for airflow through the exchanger in the case of without or with VGs. As observed in this figure, the variation of the dynamic pressure is proportional to that of the attack angle ($\alpha$). The data reported from simple baffle (of $\alpha = 90^\circ$) is also shown for comparison. Similarly to the $x$-velocity data in Figure 5 (a)–(e), the highest $Pd$ values are presented in the vicinity of the duct’s top wall, with an acceleration process which launches just following the fourth V-fin, as a consequence of the strong modifications in the flow direction in that region. The maximum dynamic pressure is obtained for $\alpha = 90^\circ$, while the lowest one is that for the cases with no VGs.

Figure 8a-g presents the distribution of the total temperature along with the exchanger for various cases, without or with different attack fins, namely $\alpha = 40^\circ, 50^\circ, 60^\circ, 70^\circ, 80^\circ,$ and $90^\circ$. As shown in the 1st case, there is no evolution in temperature gradients for both hot surfaces due to the absence of VGs. Based on the other cases; it is obvious that the temperature alters considerably through the exchanger’s hot wall for all VG cases. It can be concluded which the recirculation flows significantly affect the temperature field since superior fluid mixing can be achieved in the zones, between the heated sections and the center air areas. It leads to a great temperature gradient across the hot surface. Additionally, V-shaped baffles’ attack angle spacing ($\alpha$) influences the thermal field. Based on the above figure, it should be mentioned that temperature decreases significantly at $Re = 6 \times 10^3$ and for large $\alpha$ values. Thus, it can be concluded that the increase of the attack angle yields a decrease in the temperature of the air inside the exchanger.

Figure 9 proposes an analysis of the axial curves of velocity in the last section of the V-finned solar heat exchanger in order to assess the extent of their influence with the attack angle ($\alpha$) and flow rates ($Re$).
Figure 6. Fields of $y$-velocity for (a) smooth, and (b) 40°, (c) 50°, (d) 60°, (e) 70°, (f) 80° and (g) 90° finned heat exchangers with $Re = 6 \times 10^3$.

Figure 7. Fields of $Pd$ for (a) smooth, and (b) 40°, (c) 50°, (d) 60°, (e) 70°, (f) 80° and (g) 90° finned heat exchangers with $Re = 6 \times 10^3$.

Figure 8. Fields of $T$ for (a) smooth, and (b) 40°, (c) 50°, (d) 60°, (e) 70°, (f) 80° and (g) 90° finned heat exchangers with $Re = 6 \times 10^3$.

Next to the outlet of the exchanger, at an axial station equal to 0.829 m from the duct intake, 145 mm after the 4th VG as well as 29 × 10⁻³ m from the side right, and based on the outcomes in Figure 9a, the amount of velocity becomes around 2.00 m/s, that is 3.810 times higher than the $U_{in}$. It should be mentioned that extremely large recirculation to the back of the fourth fin led to these velocity values. The axial velocity is also presented plotted in Figure 9 versus the V-fins’ attack angle, for a constant value of $Re$ (of 6,000).

The curves indicate that the $x$-velocity improvement augments with increasing the attack value of $\alpha$; see Figure 9b. The smooth duct state is also under evaluation in this figure, which was characterized by low-value flat curves, as they were not subject to the effect of fins. The results given in Figure 9c indicate an increase of the axial velocity with the $Re$ improvement.

Figure 10 highlights the progression of the dimensionless local Nusselt number ($Nu_x/Nu_0$) measured across both lower ($y = -H/2$) and upper ($y = H/2$) exchanger walls at $\alpha = 60^\circ$ and $Re = 6 \times 10^3$. Based on the figure, the minimum value of $Nu_x$ can be situated in the lowest part of the V-fins whereas the maximum amount can be achieved on its upper sections. In the intermediate zone, considerable values can be detected because of the fluid circulation between the V-fins. The achieved results are
in perfect conformity with all researchers who declare the appearance of the local recirculation areas. And this is associated with a higher heat transfer (as an illustration, see Nasiruddin & Siddiqui, 2007). Moreover, the dimensionless numbers of Nusselt are all greater than 1, about 50 times greater than that of the exchanger with no VGs, and this reflects the effectiveness of the fins present in the duct. Additionally, the rate of air affects the heat transfer. A linear increase can be detected between $Re$ and $Nu_x$.
on both exchanger walls as listed in Figure 11a and b, respectively.

Figure 12 (a and b) reports the evolution of the dimensionless mean number of Nusselt (\(Nu/Nu_0\)) with air rate for various V-fin attack angles over the exchanger walls, \(y = -H/2\) and \(y = H/2\), respectively. It seems evident which the \(Nu/Nu_0\) is considerably increased by augmenting the rate of air since the current speed increases in the two negative and positive ways. It is also worth stating which significant enhancements can be noticed for high, low, and medium flows. The outcome can be achieved since an enhancement in the rate of the fluid increases the circulation velocity together with the convective rate of transfer of heat from the exchanger wall. Consequently, the values of the \(Nu\) augment as well. Upper \(V-\alpha\) value leads to an augmentation in the \(Nu\) values. The installation of 80° V-geometry VG performs much better than that of other 40°–70° V-fins for increased \(Nu\) number. The maximum value of \(Nu\) is around 5.969, 6.307, 6.785, 7.348 and 7.889 times larger than the smooth-walled exchanger, for \(\alpha = 40^\circ, 50^\circ, 60^\circ, 70^\circ,\) and \(80^\circ\), respectively, for \(y = -H/2\) (see Figure 12a), while around 8.419, 8.900, 9.564, 10.314 and 11.015 times for \(y = H/2\) (see Figure 12b).

However, the introduction of simple fins (of 90°) gives lower \(Nu\) values than that with \(\alpha = 40^\circ, 50^\circ, 60^\circ, 70^\circ\) and \(80^\circ\) around 3.280, 9.113, 17.369, 21.130 and 36.484%, respectively, for \(y = -H/2\) (see Figure 12a), while around 31.959, 39.493, 49.900, 61.667 and 72.648% for \(y = H/2\) (see Figure 12b), at \(Re = 3 \times 10^4\).

Figure 13 (a and b) presents the mean variations of the convective coefficient of transfer of heat (\(h_{av}\)) with the air rate at different attacks (\(\alpha\)) over the exchanger walls, top and lower exchange surface, respectively. As expected, the obtained \(h_{av}\) employing the V-geometry fin with \(\alpha = 80^\circ\) is substantially upper than that with smaller \(\alpha\) values. At \(Re = 3 \times 10^4\) and \(y = -H/2\), the case of \(\alpha = 80^\circ\) provides a value of \(h_{av}\) that is higher by about 22.142, 16.197,
11.376 and 5.418% than those given with \( \alpha = 40°, 50°, 60°, \) and 70°, respectively (see Figure 13a). However, and at \( y = H/2 \), the superiority of the case \( \alpha = 80° \) in terms of \( h_{av} \) was about 22.221, 16.509, 11.219 and 5.356% over those reached by \( \alpha = 40°, 50°, 60°, \) and 70°, respectively (see Figure 13b). It was also observed that the augmentation of Reynolds number yields an increase in \( h_{av} \) values for all \( \alpha \) cases.

On the other hand, the presence of a simple fin, of 90° attack, reduces the \( h_{av} \) values by about 3.862, 11.794, 18.225, 26.173 and 33.401%; and 34.200, 44.056, 53.184, 63.301 and 72.543, compared to \( \alpha = 40°, 50°, 60°, 70° \) and 80° cases, at \( y = -H/2 \) and \( y = H/2 \), respectively.

The regional variation of the coefficient of friction, \( C_f \) over \( y = -H/2 \) for V-geometry fins’ attack angle values of \( 40° \leq \alpha \leq 80° \) and \( Re = 6,000 \) is illustrated below in Figure 14a. The smallest \( C_f \) is found around the lower surface-installed V-VGs, while the largest value is encountered near the area confronting the upper wall-fixed V-fins. Moreover, the value of \( C_f \) augments for augmenting the V-angles. The effect of a simple VG on the \( C_f \) curves is also evident in Figure 14a, as it shows an increase in the values of friction across the lower gaps, due to the presence of the upper fins, with average values generated by the recycling cells. The skin friction coefficient is also influenced by the rate of air. It is clear that the \( C_f \) becomes more substantial once the Reynolds number augments; see Figure 14b.

The axial change of \( C_f \) curves for the duct’s hot top axis with distinct flow attack values at \( Re = 6,000 \) is graphed in Figure 15a. It seems obvious that the greatest amounts are initially adjacent to the duct’s outlet since the fourth V-fin directs the flow towards the exchanger’s upper part with great velocities and then in the intermediate zones since the fluid downstream of the upper wall-placed V-VG recirculates. Furthermore, the attack angle of the

![Figure 14.](image1.png)  
![Figure 15.](image2.png)

**Figure 14.** \( C_f \) over the lower exchanger wall for various situations.  
**Figure 15.** \( C_f \) over the upper exchanger wall for various situations.
V-fin also affects the $C_f$ value. A direct relationship is noted between $C_f$ and $\alpha$. The lowest $C_f$ value was found at $\alpha = 40^\circ$, at all locations. This value augments as $\alpha$ goes up. The friction curves caused by the behavior of simple fins are also included in Figure 15a according to the same $Re$ number used. A qualitative behavior similar to that reported by V-fins is recorded. Also, Figure 15b illustrates the impact of the flow rate on the axial progress of the skin friction coefficient $C_f$, with respect to Reynolds numbers. Identical results can be observed by this figure indicating $C_f$ value is augmented by increasing the rate of air.

The curves are shown in Figure 16 stand out the progress of the average coefficient of friction, $f$ with respect to the Reynolds number (6,000–30,000), at five various V-baffles’ attack angle values ($\alpha = 40^\circ, 50^\circ, 60^\circ, 70^\circ$, and $80^\circ$). Generally, the $f$ has a tendency to augment with increasing $Re$. It can be noted that the improvement of the $f$ factor is more significant than that of the $Nu$ and $h_{av}$ parameters. This might recommend which the turbulence fields and velocity enlarge faster than the isotherms. Based on the graph, it can be observed that the $f$ factor augments once the V-geometry obstacles’ attack angle ($\alpha$) increases, independently from the Reynolds numbers.

**Figure 19.** Hydrothermal fields for various 40°V-fins’ heights, $D_t = 0.412$ m and $Re = 6,000$. 
number. For an air rate of \( Re = 3 \times 10^4 \) and compared with the simple finned duct, of 90°, the insertion of V-fins with \( \alpha = 70° \), and 80° provided an elevation in the \( f \) factor, by about 9.354 and 35.634 percent, respectively, at \( y = -H/2 \) (see Figure 16a). However, at \( y = H/2 \) and for the same attack angles, the augmentation of the \( f \) factor was about 15.246 and 42.256 percent over the 90° finned duct (see Figure 16b). On the other hand, and at the same upper air rate, of 30,000, the superiority of the case \( \alpha = 90° \) in terms of friction was about 44.988%, 32.516% and 14.476%; and 41.056%, 28.571% and 10.084% over those reached by \( \alpha = 40° \), 50°, 60°, and 70°, at \( y = -H/2 \) and \( y = H/2 \), respectively.

Figure 17 assesses the performance (TEF) of all proposed finned duct heat exchangers for different \( \alpha \) attacks from a specific range of rate of air. The TEF exceeded the value 1 in all considered cases, i.e. there is a hydrothermal improvement compared to the duct with smooth walls. In the case of large flow rates, the TEF drop for V-finned ducts with smaller attack angle is determined to be greater than that with a larger attack angle. The TEF decreases by the increment in V-shaped baffles’ attack angle (\( \alpha \)) and thus, the \( \alpha = 40° \) provides maximum TEF.

![Figure 16. Effect of rate of air on \( f \) over the exchanger walls for different attack V-fins.](image1)

![Figure 17. Performance evaluation for different V-finned heat exchangers.](image2)

![Figure 18. TEF comparison with other VG models.](image3)
The 50°, 60°, 70°, 80° and 90° attack fins present diminutions of about 1.323%, 1.775%, 1.936%, 1.957% and 20.577%, respectively, for $y = -H/2$ and $Re = 30,000$ (see Figure 17a), while about 0.744%, 1.159%, 1.770%, 2.158% and 11.267% for $y = H/2$ (see Figure 17b), in $TEF$ relative to that in the 40° $V$-shaped fin indicating the 40° V-type baffle is more beneficial than the other cases.

On the other hand, the current optimum V-fin of 40° was also evaluated along with other previous obstacle models from the literature as reported in Figure 18. For $Re = 1.2 \times 10^4$ and $y = H/2$, this fin exhibits an increase in $TEF$ values of about 1.196%, 13.528%, 21.340%, 19.447%, 23.113%, 14.462%, 17.572%, 23.779% and 22.357% over that of cases with, four flat (Menni, Azzi, and Zidani, 2018); flat and upstream V (Menni, Azzi, Didi, et al., 2018), flat and trapezoidal (Menni & Azzi, 2018), flat and triangular (Menni & Azzi, 2018), flat and Z (Menni & Azzi, 2018), flat and W (Menni & Azzi, 2018), flat and down V (Menni et al., 2019), two rectangular (Menni et al., 2020), and three S- VGs (Chamkha et al., 2020), respectively.

So, the current model according to the attack angle of 40° is considered as an ideal solution for determining the finned duct solar air-heat exchangers, especially in the case of high flow rates.

### Effect of the 40° V-fin height

In this section, the effects of the V-fin height on the overall characteristics of the exchanger are highlighted for constant V-attack and space values ($\alpha = 40°$ and $D_s = 0.142$ m). For this purpose, five geometrical configurations were realized, namely: 0.04, 0.06, 0.08, 0.10, and 0.12 m, respectively. Theses selected values of the V- height correspond to $H_b/2$, $3H_b/4$, $H_b$, $5H_b/4$, and $3H_b/2$, respectively.

**Figure 20.** Effect of rate of air on $Nu/Nu_0$ over the duct walls for different height 40°V-fins.

**Figure 21.** Effect of rate of air on $f$ over the exchanger walls for different height 40°V-fins.
The flow patterns induced for each geometrical case are illustrated in Figure 19a. It seems that the recirculation loops generated before and after the V-fin are enlarged with respect to the increasing value of V-length. The inclination of the segments of the V-fin by 40° presents a power generator of vortices, giving thus a strong interaction between the fluid particles. The mean velocity is also intensified with the rise, where the maximum value \(V_{max}\) passed at \(Re = 6,000\) from 1.106 to 3.577 m/s when the V-section height changed from 0.04 to 0.12 m, i.e. from \(H_b/2\) to \(3H_b/2\) (see Figure 19b). Furthermore, the excessive amount of V-height yields regions of strong velocities at the tip of each segment of the fin, with moderate or low velocities elsewhere.

These regions, where the high velocities occur, are also the location of significant values of the dynamic pressure (see Figure 19c). The maximum value of the pressure \(Pd_{max}\) at \(Re = 6,000\) passed from 0.749 Pa to 8.260 Pa when the V-length changed from 0.04 to 0.12 m, which corresponds to an increase by about 11 times compared with the case \(H_b = 0.04\) m. This is mainly due to the flow behavior generated by this kind of vortex generators.

Certainly, the flow characteristics have a great impact on thermal behavior. From the plots provided in Figure 19d, the thermal fields are uniform until reaching the 1st V-fin, where the thermal exchange begins clearly. The heat exchange is then intensified gradually when the fluid passes through the next fins. Consequently, the increasing V-height reduces the required length in the exchanger to obtain the desired temperature of the fluid, which results in a compact exchanger.

Further details on the thermal exchanger are given in Figure 20a and b, where the values of the normalized Nu number, \(Nu/Nu_0\), are presented for the lower and upper hot walls of the duct, respectively. The variations of \(Nu/Nu_0\) are presented for different values of Reynolds number and 40° V-fin height. The profiles of both figures confirm the finding illustrated in Figure 19d. The \(Nu/Nu_0\) increases with the rise of the fin height. More precisely and for sizes, 0.04 and 0.12 m, respectively, the most significant values of \(Nu/Nu_0\) are 5.028 and 7.983 on the lower wall of the duct, while 6.599 and 10.927 on the upper wall. These values are obtained at \(Re = 3 \times 10^4\). Compared with the case \(H_b = 0.04\) m, the normalized \(Nu\) number has been increased by about 1.587 and 1.655 times for the lower and upper walls of the duct, respectively.

In addition, the increase in Reynolds number produces an augmentation of the values of Nusselt number due to the intensification of the interaction between the fluid particles that is resulted in high \(Re\). However, the friction factor is also increased on both walls of the exchanger with the rise of \(Re\), as observed in Figure 21a and b. The augmentation of the V-fin height has proved its efficiency in terms of enhancement of the thermal exchange, but the friction factor is also increased (see Figure 21). At \(Re = 30,000\), the value of the friction factor passes from 0.017 to 0.076 on the lower wall, and from 0.031 to 0.141 on the upper wall, when \(H_b\) changed from \(H_b/2\) to \(3H_b/2\).

Finally, and aiming to give a global evaluation of the efficiency of the exchanger, the thermal enhancement factor (TEF) is determined for different flow rates and geometrical conditions (see Figure 22). For both walls of the duct, it seems that the TEF increases continually with the rise of the flow rates (Re). However, and when changing the V-fin height, the TEF increases until \(H_b = 0.08\) m, then it decreases again. At \(H_b = 0.08\), the values of TEF for the upper and lower walls of the exchanger are 2.163 and 1.926, respectively. The comparison between the values of the TEF on the lower and upper walls reveals a slight difference. It is slightly elevated in the case of \(y = H/2\). This is maybe due to the strong gaps in the

![Figure 22](image-url). Performance evaluation for various height 40° V-fins.
upper exchanger side, in contact with the top horizontal axis, where they have high flow speeds, due to the high dynamic-pressure which produces high temperature gradients in the region and thus, an important heat transfer, in addition to the intensity of the recycling cells localized behind the upper V-fins.

**Effect of the space between 40° V-fins**

Another geometrical parameter that may affect the hydrothermal characteristics of the exchanger is studied in the section. It concerns the spacing between 40° V-fins (or the distance of separation). For this objective, four cases were considered, which are: 0.071, 0.1065, 0.142, and 0.1775 m. These values correspond to $D_s/2$, $3D_s/4$, $D_s$, and $5D_s/4$, respectively.

At $Re = 6,000$ and for an exchanger equipped with 40° V-fins having $H_b = 0.08$ m, the hydrodynamics induced for the four cases are plotted in Figure 23a under a two-dimensional view. As observed, the location of the vortices is the same in each case. However, the intensity of these recirculation loops increases with the reduced distance of separation between the baffles. In the first case ($D_s/2$), the shape of the vortex developed after the first fin is highly affected by the second fin due to the small spacing between the VGs. The maximum value of the average velocity in this case ($D_s/2$) is the highest compared with the three other cases (see Figure 23b). It reaches 5.126 m/s at $D_s/2$, while it is only around 1.689 m/s at $5D_s/4$.

The dynamic-pressure is also strongly affected when changing the distance of separation between the fins (see Figure 23c). The most significant amount of the dynamic-pressure that is around 17.558 Pa is obtained with the 1st case ($D_s/2$), which is higher by about 10.038 times than the maximum value given by the last case ($5D_s/4$). This great difference in the $Pd$ between the first and last cases

![Figure 23](image-url)
is resulted by the wall effect, which increases furthermore with the reduced space between the VGs.

The intensification of the interaction between the fluid particles that is resulted by the small distance of separation promotes the heat transfer rates. As observed in Figure 23d, the cold fluid is rapidly heated during its passage through the finned region of the exchanger when the distance \( D_s \) is small. Therefore, the case \( D_s/2 \) may be selected as the most efficient in terms of enhanced thermal exchange. In addition, the increasing space between the vortex generators participates in the augmentation of the length in the duct that is required to obtain the desired temperature of the fluid.

Further insight into the convection phenomenon that is generated when changing the flow rate (Reynolds number), and the V-space (spacing) is provided in Figure 24a and b. In these figures, the normalized values of Nusselt number \( Nu/Nu_0 \) are given for the lower and upper walls of the exchanger, respectively. Increasing values of the \( Nu/Nu_0 \) are observed with the rise of \( Re \). However, an inverse proportionality is remarked when increasing the V-fins’ distance. However, similar behavior of the variation of \( Nu/Nu_0 \) is shown for both walls of the exchanger. Overall, the most significant values of \( Nu/Nu_0 \) are reached at the highest \( Re \) (\( Re = 30,000 \) here) and \( D_s/2 \), which are 6.661 and 9.111 on the lower and upper walls, respectively. While the lowest values of \( Nu/Nu_0 \) at the same \( Re \) (30,000) are those of \( 5D_s/4 \), which are 5.755 and 8.204 on the lower and upper walls, respectively.

The reduced distance between 40° V-fins has been found to be advantageous in terms of improvement of the heat transfer. However, the friction factor \( (f) \) has also been increased accordingly. The variation of \( f \) vs. \( D_s \) and \( Re \) is highlighted in Figure 25. At \( Re = 30,000 \), the highest amount of \( f \) is reached with \( D_s/2 \), and it is equal to 0.0379 and 0.0624 on the lower and upper walls, respectively. At the same \( Re \), the lowest values of \( f \) are 0.0185 and 0.0430 on the lower and upper walls, respectively. The analysis of these results reveals an increase in the friction factor by about 204.864% and 145.116% on the lower and upper walls.

**Figure 24.** Effect of rate of air on \( Nu/Nu_0 \) over the duct walls for various 40°V-fins’ spaces.

**Figure 25.** Effect of rate of air on \( f \) over the exchanger walls for various 40°V-fins’ spaces.
upper walls, respectively, when the space between V-fin is changing from \(5D_s/4\) to \(D_s/2\).

**Conclusions**

In this research work, it is focused mainly on important geometric variables related to V-finned duct solar air-heating exchangers, namely the attack angle (40°–90°), the height \((H_b/2–3H_b/2)\) as well as the separation length between VGs \((D_s/2–5D_s/4)\), along with the flow rate \((6,000–30,000)\), to reach an optimal solution for hydrodynamic behavior and heat transfer performance.

The V-fin’s attack angle impacts the flow and heat transfer characteristics. If the attack angle \((\alpha)\) is increased, it leads to the flow deviation and acceleration adjacent to the V-fins, which results in intensifying the convective rate of heat transfer. While this attack angle augmentation is restricted since it might cause the pressure loss to augment. It seems evident which more sensible effects on the upper left side of the V-fin can be seen by alterning the V-angle \((\alpha)\) because of the air-field direction deviation. On the other hand, increasing the \(\alpha\) angle from 80° to 90° increases the field intensity and enhances the recirculation cells, giving maximum values of \(x\)- and \(y\)-velocities. Additionally, V-shaped fins’ attack angle spacing \((\alpha)\) influences the thermal field. Based on the considered attack values, it should be mentioned that temperature decreases significantly for large \(\alpha\) values. Thus, it can be concluded that the increase of the attack angle yields a decrease in the temperature of the air inside the exchanger. The heat transfer in terms of \(\text{Nu}\) number is around 5.969, 6.307, 6.785, 7.348 and 7.889 times larger than the smooth-walled exchanger, for \(\alpha = 40°, 50°, 60°, 70°,\) and 80°, respectively, for \(y = -H/2\), while around 8.419, 8.900, 9.564, 10.314 and 11.015 times for \(y = H/2\). However, the introduction of simple fins, of 90°, gives lower \(\text{Nu}\) values than that with \(\alpha = 40°, 50°, 60°, 70°\) and 80° around 3.280, 9.113, 17.386, 27.130 and 36.484%, respectively, for \(y = -H/2\), while around 31.959, 39.493, 49.900, 61.667 and 72.648% for \(y = H/2\), at \(Re = 3 \times 10^4\). Furthermore, the attack angle of the V-fin also affects the \(f\) value. A direct relationship is noted between \(f\) and \(\alpha\). The lowest \(f\) value was found at \(\alpha = 40°\), at all locations. For a maximum air rate and compared with the simple finned duct, of 90°, the insertion of V-fins with \(\alpha = 70°\), and 80° provided an elevation in the \(f\) factor, by about 9.354 and 35.634 percent, respectively, at \(y = -H/2\). However, at \(y = H/2\) and for the same attack angles, the augmentation of the \(f\) factor was about 15.246 and 42.256 percent over the 90° finned duct. On the other hand, and at the same upper air rate, of 30,000, the superiority of the case \(\alpha = 90°\) in terms of friction was about 44.988%, 32.516% and 14.476%; and 41.056%, 28.571% and 10.084% over those reached by \(\alpha = 40°, 50°, 60°,\) and 70°, at \(y = -H/2\) and \(y = H/2\), respectively. In terms of the performance, The \(\text{TEF}\) decreases by the increment in V-shaped baffles’ attack angle \((\alpha)\) and thus, the \(\alpha = 40°\) provides maximum \(\text{TEF}\). The 50°, 60°, 70°, 80° and 90° attack fins present diminutions of about 1.323%, 1.775%, 1.936%, 1.957% and 20.577%, respectively, for \(y = -H/2\) and \(Re = 30,000\), while about 0.744%, 1.159%, 1.770%, 2.158% and 11.267% for \(y = H/2\), in \(\text{TEF}\) relative to that in the 40° V-shaped fin indicating the 40° V-type baffle is more beneficial than the other cases.

On the other hand, improving the length of the V-fins or decreasing in the space between them, increases the flow strength by enlarging the recycling cells, which reflects on the hydrodynamic behavior and changes the heat transfer. The presence of this new model of fins also highlights a hydrothermal improvement ranging between 1.196 and 23.779 percent compared to the previously mentioned models, reflecting the effectiveness of the new system of solar heat exchangers with air V-finned ducts. So, the current model according to the attack angle of 40° is considered as an ideal solution for determining the finned duct solar air-heat exchangers, especially in the case of high flow rates.

For future studies, the introduction of pores in the V-geometry baffles may provide further reduction in the pressure drop. The newly suggested design of VGs may also participate in the elimination of the low heat transfer areas (LHTA) that form behind the baffles. It is also proposed to examine the performance of the optimized geometrical configuration presented in our study (40° V-baffle) for other industrial applications, such as the cooling of complex non-Newtonian fluids.

**Disclosure statement**

No potential conflict of interest was reported by the author(s).

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